

# DPMS: Advanced Materials

## Lecture Y4: Computational Optics I

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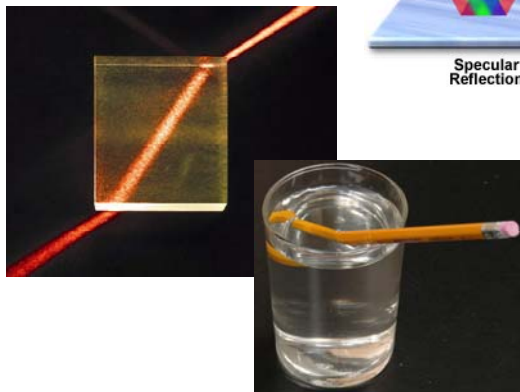
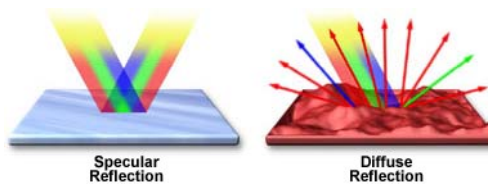
# Photonic and optoelectronic technologies

- Understanding light-matter interactions is crucial for new applications



# Reflection and refraction

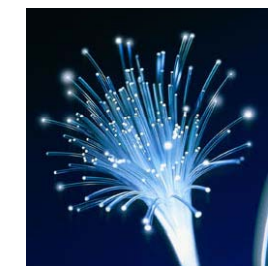
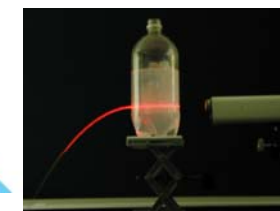
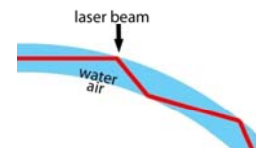
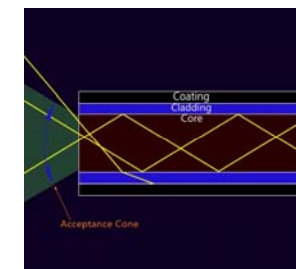
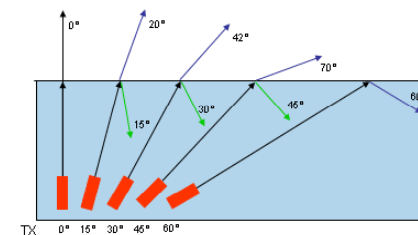
- Material property:
  - dielectric function
  - index of refraction



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# Trapping and guiding light

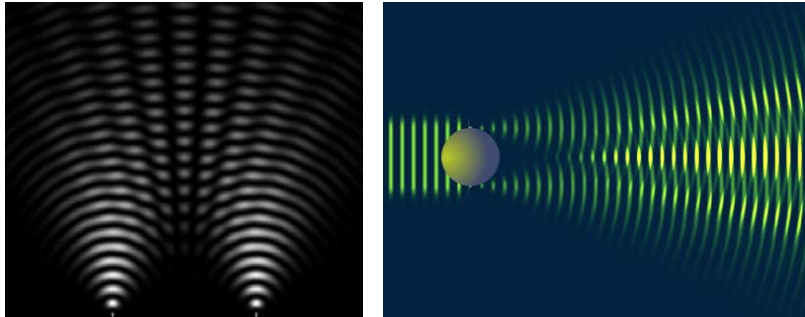
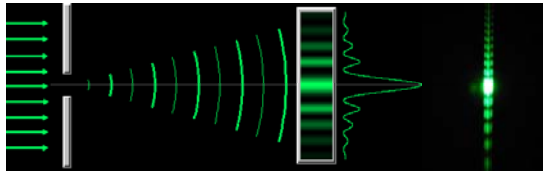
- Total internal reflection



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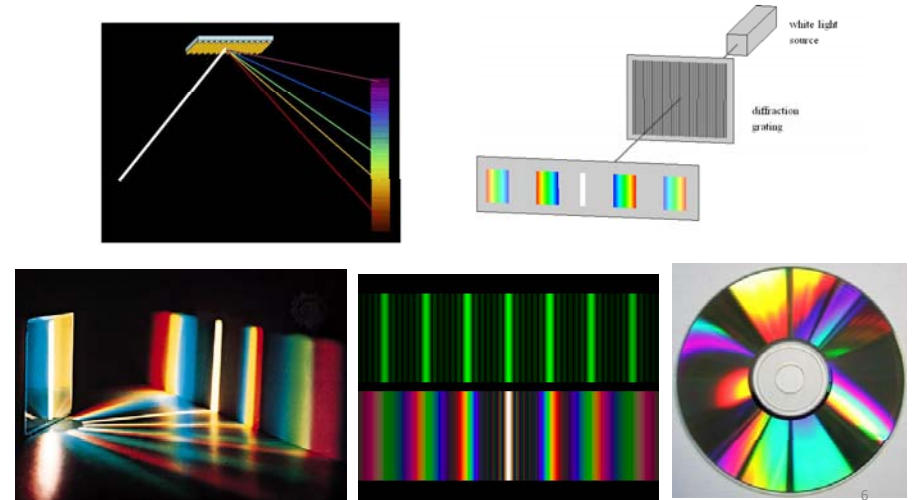
## Diffraction and interference

- Diffracted waves interfere



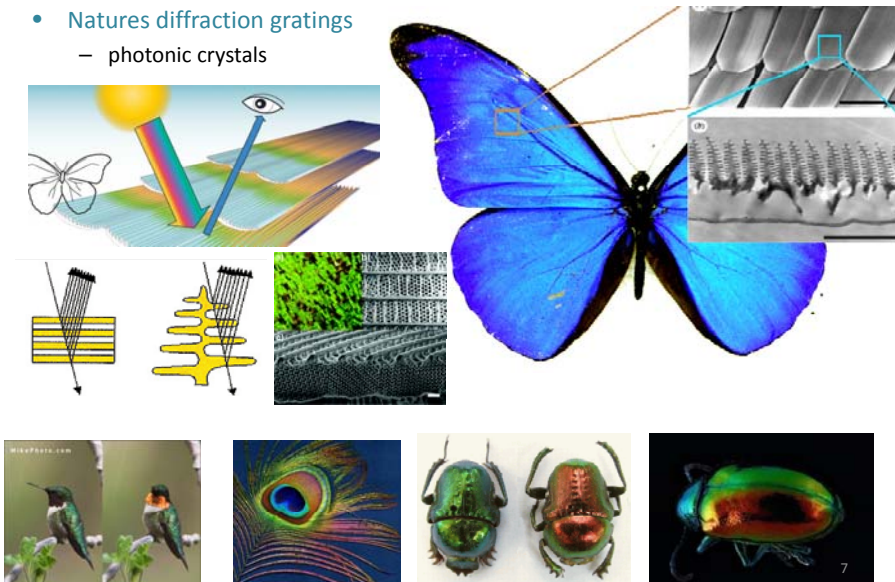
## Diffraction gratings

- Diffraction gratings are used for wavelength separation



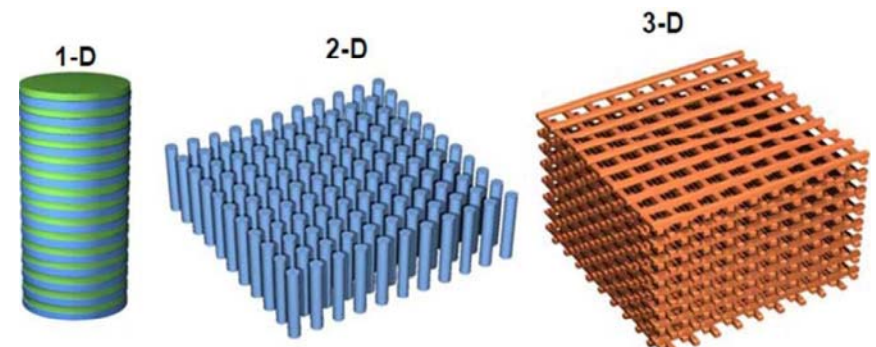
## Structural color

- Nature's diffraction gratings  
– photonic crystals



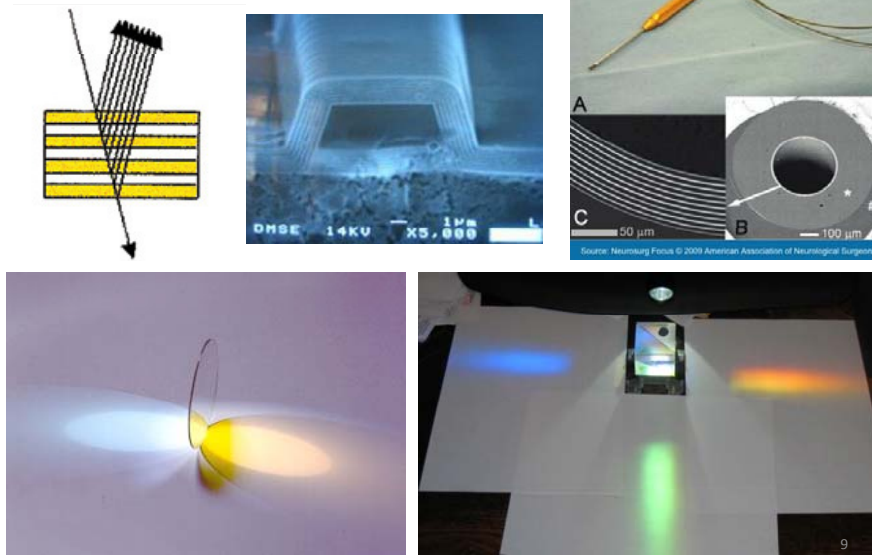
## Photonic crystals in technology

- Applications in extreme light manipulation

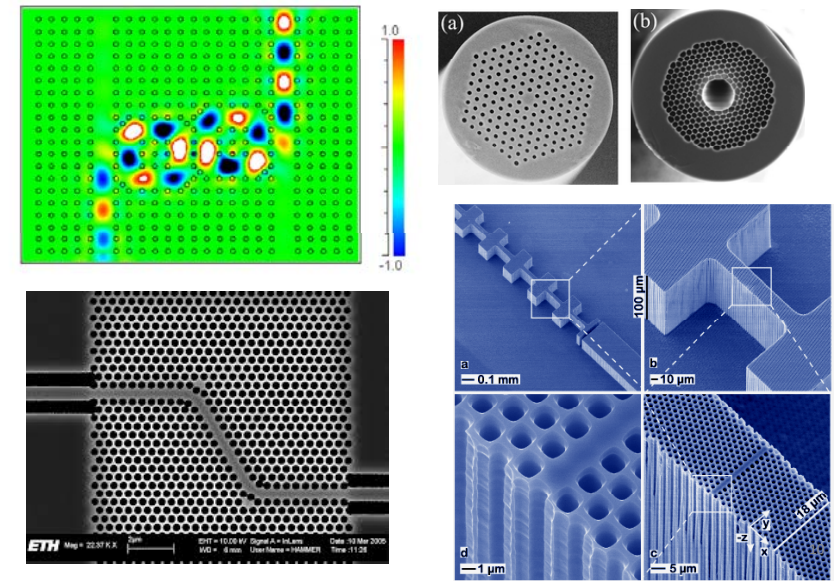




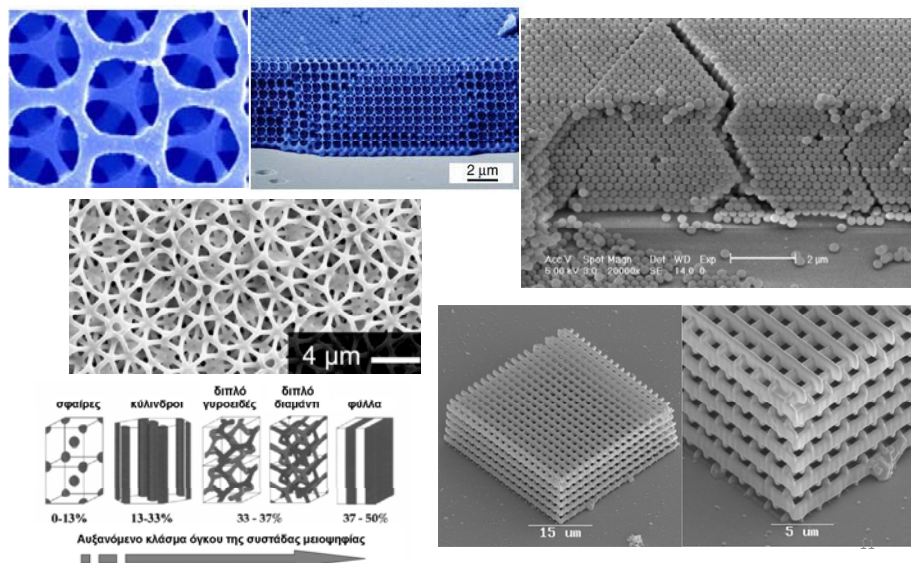
## Examples of 1D photonic crystals



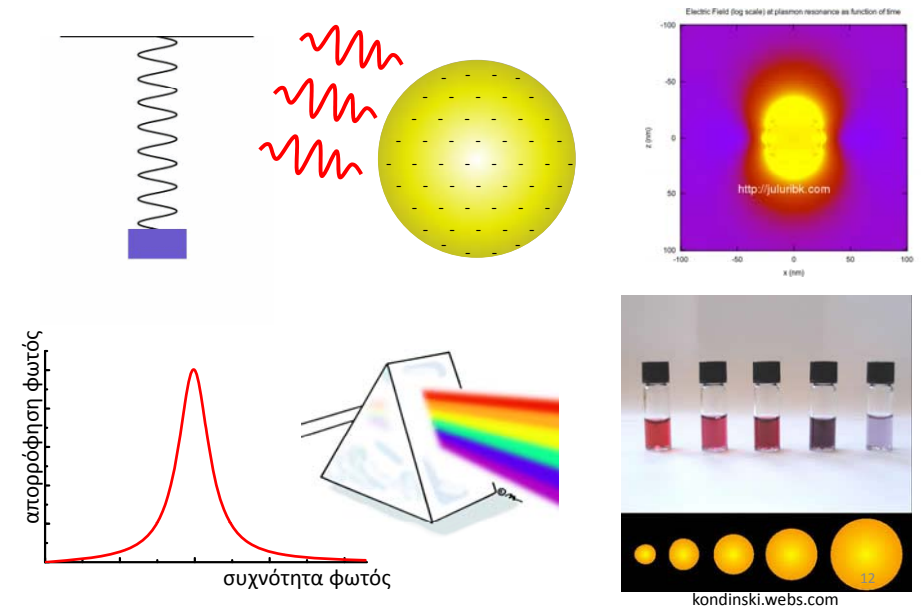
## Examples of 2D photonic crystals



## Examples of 3D photonic crystals

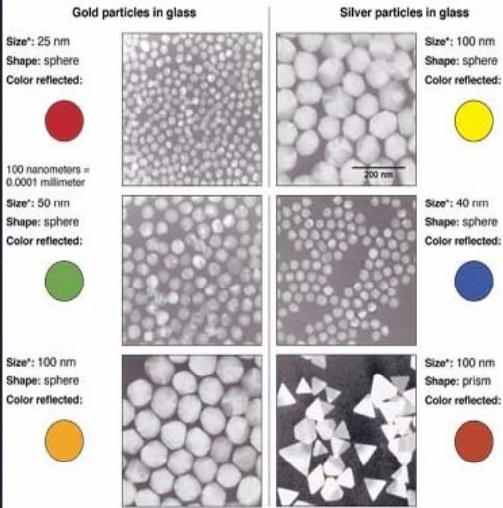


## Plasmonics: interactions with nanostructured metals





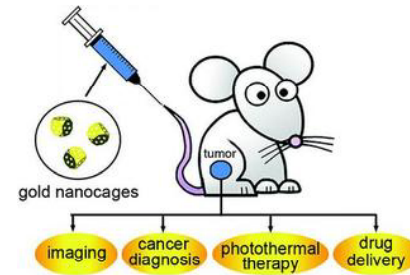
## Plasmonics in the middle ages



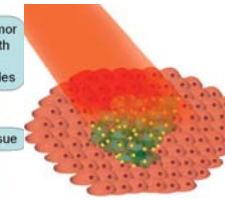
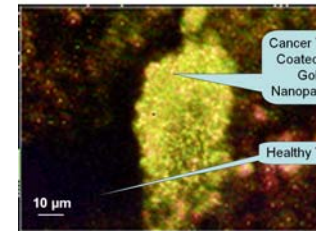
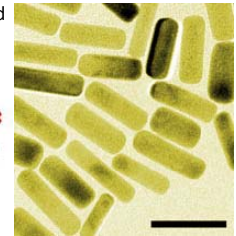
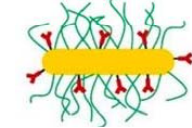
Institute of Nanotechnology, Northwestern University



## Εφαρμογές πλασμονικών νανοσωματιδίων στην ιατρική



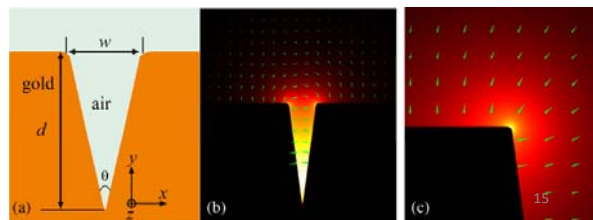
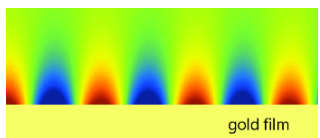
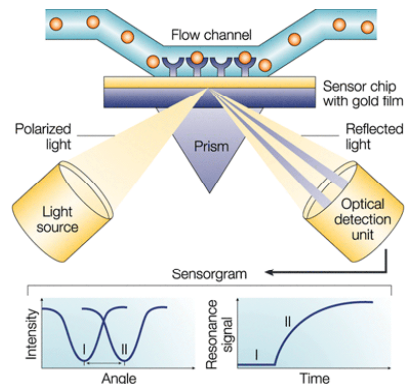
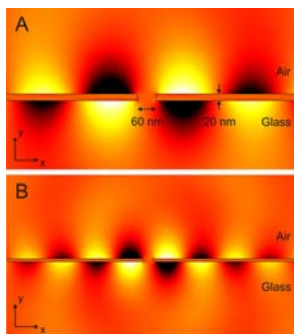
anti-EGFR conjugated gold nanoparticles



Rice University



## Surface plasmons



## Outline

- E/M waves
- Wave equations
- Boundary conditions
- Materials' optical response
- Materials' optical characterization
- Numerical solutions
- Applications design



## The classical nature of light

- Oscillating electric and magnetic fields

- propagation speed  $c=3 \times 10^8$  m/s

- Transverse waves

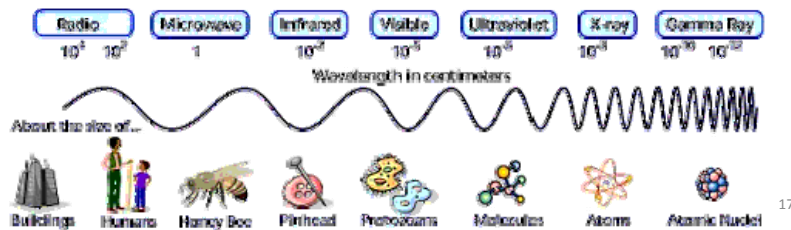
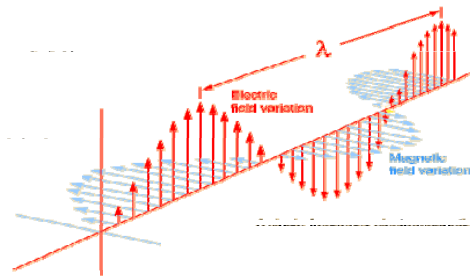
- perpendicular to propagation direction
- perpendicular to each other

- Gets polarized

- linearly polarized (e.g. x and y)
- circularly polarized (e.g. right and left)
- generally two orthogonal polarizations

- Interaction with materials

- reflection, refraction, absorption, interference, scattering, diffraction, ...



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## E/M wave properties

- Wavelength  $\lambda$

- spatial periodicity

- Period  $T$

- temporal periodicity

- Wavevector  $k = 2\pi / \lambda$

- spatial repetition rate

- Angular frequency  $\omega = 2\pi / T$

- temporal repetition rate

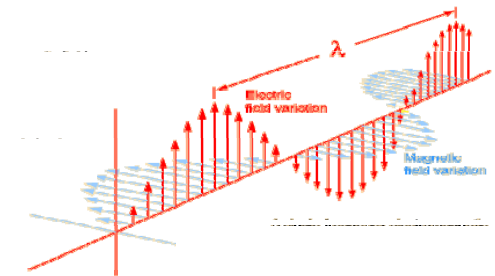
- Dispersion relation  $k = n\omega / c$  ( $n$  is index of refraction)

- relation between  $k$  and  $\omega$
- relation between space and time
- relation between momentum and energy ( $E = \hbar\omega$ ,  $p = \hbar k$ )

- In 3D space,  $k$  is a 3D vector

- Electric and magnetic fields are 3D vectors

- polarization of wave is the electric field polarization



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## E/M wave equations

- Maxwell's equations

$$\nabla \cdot \mathbf{D} = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = +\frac{\partial \mathbf{D}}{\partial t}$$

$$\mathbf{D} \equiv \mathbf{D}(\mathbf{r}, t)$$

$$\mathbf{B} \equiv \mathbf{B}(\mathbf{r}, t)$$

$$\mathbf{E} \equiv \mathbf{E}(\mathbf{r}, t)$$

$$\mathbf{H} \equiv \mathbf{H}(\mathbf{r}, t)$$

- Constitutive relations

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon_0 \epsilon_r \mathbf{E} = \epsilon \cdot \mathbf{E}$$

$$\mathbf{B} = \mu_0 (\mathbf{H} + \mathbf{M}) = \mu_0 \mu_r \mathbf{H} = \mu \cdot \mathbf{H}$$

- Material properties

- electric permittivity

$$\epsilon \equiv \epsilon(\mathbf{r})$$

- magnetic permeability

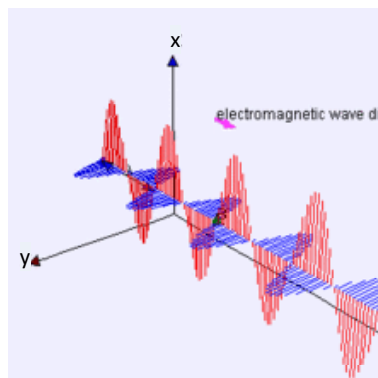
$$\mu \equiv \mu(\mathbf{r})$$

- refractive index

$$n = n(\mathbf{r}) = \sqrt{\epsilon_r(\mathbf{r}) \mu_r(\mathbf{r})}$$

- At optical frequencies

$$\mu_r = 1$$



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## E/M wave equations in 1D

- Maxwell's equations

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \Rightarrow \frac{\partial E}{\partial z} = -\mu_0 \frac{\partial H}{\partial t}$$

$$\nabla \times \mathbf{H} = +\frac{\partial \mathbf{D}}{\partial t} \Rightarrow -\frac{\partial H}{\partial z} = \epsilon_0 \epsilon_r \frac{\partial E}{\partial t}$$

- Combine for one unique equation

$$\frac{\partial E}{\partial z} + \mu_0 \frac{\partial H}{\partial t} = 0 \Rightarrow \frac{\partial}{\partial z} \left( \frac{\partial E}{\partial z} + \mu_0 \frac{\partial H}{\partial t} \right) = 0$$

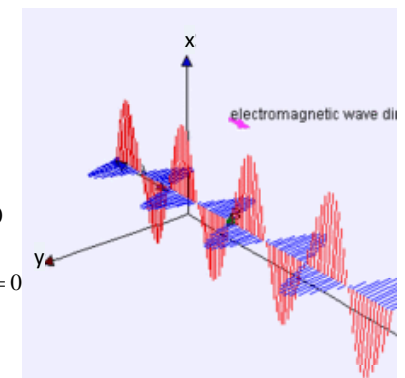
$$\Rightarrow \frac{\partial^2 E}{\partial z^2} + \mu_0 \frac{\partial}{\partial t} \frac{\partial H}{\partial z} = 0 \Rightarrow \frac{\partial^2 E}{\partial z^2} - \epsilon_0 \mu_0 \epsilon_r \frac{\partial^2 E}{\partial t^2} = 0$$

$$\Rightarrow \frac{\partial^2 E}{\partial z^2} - \frac{\epsilon_r}{c^2} \frac{\partial^2 E}{\partial t^2} = 0 \quad c^2 = 1/(\epsilon_0 \mu_0)$$

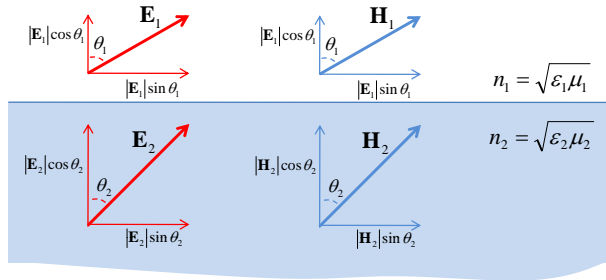
- Wave solution are plane waves

$$\mathbf{E} = \hat{x} E_0 e^{i(kz - \omega t)}$$

$$-k^2 E_0 e^{i(kz - \omega t)} + \frac{\epsilon_r \omega^2}{c^2} E_0 e^{i(kz - \omega t)} = 0 \Rightarrow k = \sqrt{\epsilon_r} \omega / c = n\omega / c_0$$



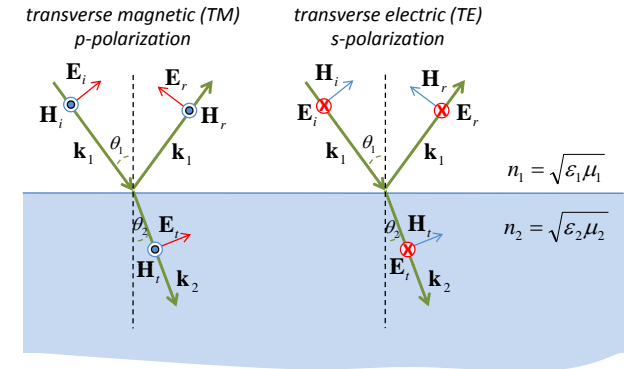
## Boundary conditions at interfaces



### Continuity at the boundary:

- of the parallel components of the electric field  $\mathbf{E}_{//}$   $|\mathbf{E}_1| \sin \theta_1 = |\mathbf{E}_2| \sin \theta_2$
- of the perpendicular components of the displacement  $\mathbf{D}_{\perp}$   $\epsilon_1 |\mathbf{E}_1| \cos \theta_1 = \epsilon_2 |\mathbf{E}_2| \cos \theta_2$
- of the parallel components of the magnetic field  $\mathbf{H}_{//}$   $|\mathbf{H}_1| \sin \theta_1 = |\mathbf{H}_2| \sin \theta_2$
- of the perpendicular components of the inductance  $\mathbf{B}_{\perp}$   $\mu_1 |\mathbf{H}_1| \cos \theta_1 = \mu_2 |\mathbf{H}_2| \cos \theta_2$

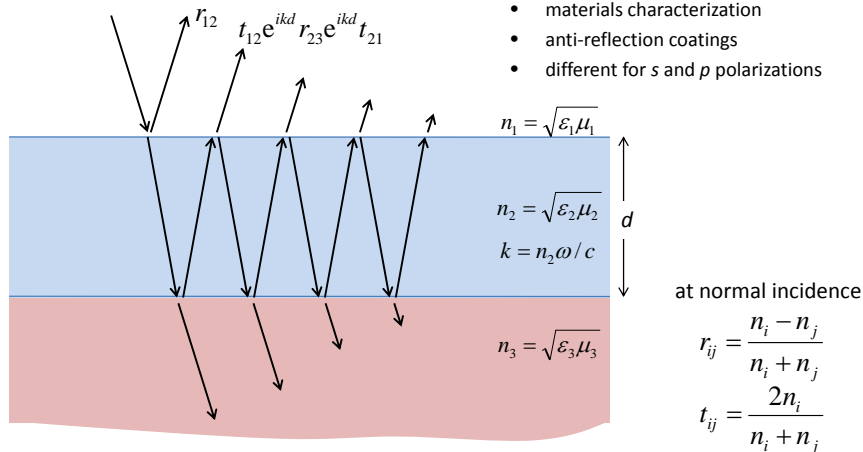
## Reflection and refraction at an interface



- Boundary conditions on  $\mathbf{k}$** 
  - parallel component of  $\mathbf{k}$  conserved  $|\mathbf{k}_1| \sin \theta_1 = |\mathbf{k}_2| \sin \theta_2$   $|\mathbf{k}| \equiv k = n\omega/c$
  - momentum conservation  $n_1 \sin \theta_1 = n_2 \sin \theta_2$  Snell's law
- Boundary conditions on fields**
  - reflection and transmission amplitudes  $r_{12}, t_{12}$
  - reflectivity and transmittance  $R = |r_{12}|^2, T = (n_2/n_1) |t_{12}|^2$
- Inside materials**  $\lambda = \lambda_0/n, v = c/n$

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## Reflection and transmission from a thin film



### Total reflection

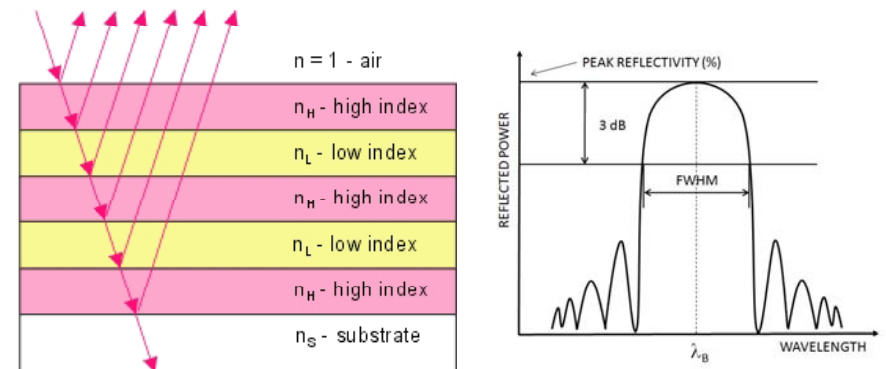
$$r = r_{12} + t_{12} e^{ikd} r_{23} e^{ikd} t_{21} + t_{12} e^{ikd} r_{23} e^{ikd} r_{21} e^{ikd} r_{23} e^{ikd} t_{12} + \dots$$

$$= r_{12} + t_{12} r_{23} t_{21} e^{i2kd} \cdot [1 + r_{21} r_{23} e^{i2kd} + (r_{21} r_{23} e^{i2kd})^2 + \dots]$$

$$= r_{12} + \frac{t_{12} r_{23} t_{21} e^{i2kd}}{1 - r_{21} r_{23} e^{i2kd}}$$

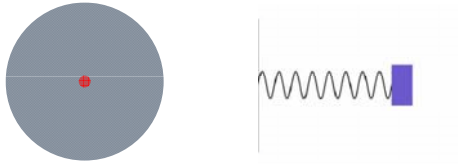
## Reflection from periodic film arrays

- 1D photonic crystal – Bragg mirror**
  - multiple reflections result into extreme reflection for some frequencies
  - perfect mirrors for lasers, selective filters, sensors, etc



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# Dielectric function: material polarizability



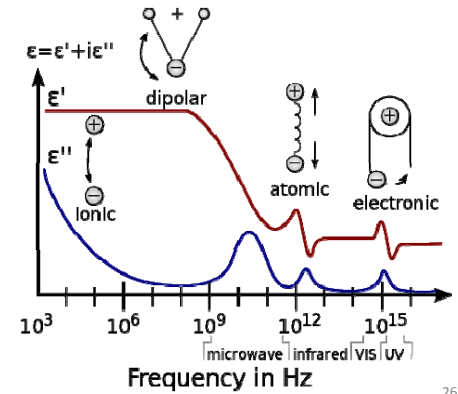
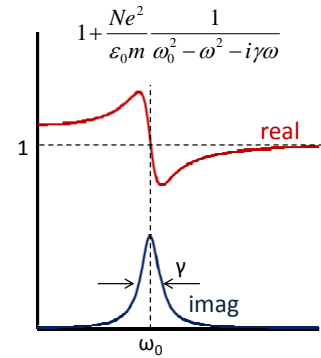
- Driven oscillation  $m \frac{d^2r}{dt^2} = -m\gamma \frac{dr}{dt} - Kr - eE$   $\omega_0 = \sqrt{K/m} \Rightarrow K = m\omega_0^2$
- Assume harmonic oscillation  $r = r_0 e^{-i\omega t}$   $E = E_0 e^{-i\omega t}$   
 $-\omega^2 m r = i\omega m \gamma r - m\omega_0^2 r - eE \Rightarrow r = -\frac{eE}{m} \frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega}$
- Polarization depends on number density of electronic orbitals  
 $P = -Ner = \frac{Ne^2}{m} \frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega} E$   $D = \epsilon_0 \epsilon_r E = \epsilon_0 E + P = \epsilon_0 \left( 1 + \frac{Ne^2}{\epsilon_0 m} \frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega} \right) E$
- Dielectric function  $\epsilon_r = 1 + \frac{Ne^2}{\epsilon_0 m} \frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega}$

# Dielectric function

- Each orbital has a different contribution

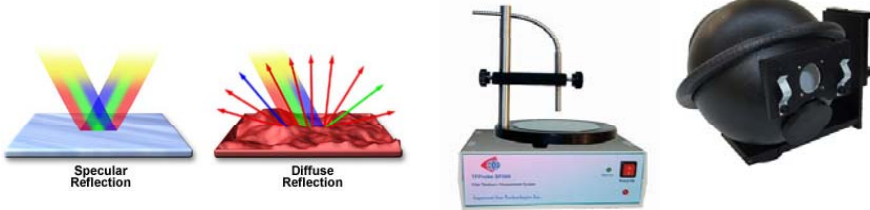
- electronic: UV and Visible
- atomic: IR
- etc

$$\epsilon_r = 1 + \sum_i \frac{N_i e^2}{\epsilon_0 m_i} \frac{1}{\omega_{0i}^2 - \omega^2 - i\gamma_i \omega}$$

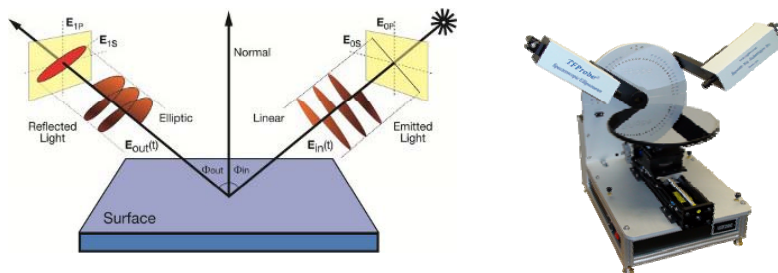


# Materials' optical characterization

- Reflectance spectroscopy



- Ellipsometry



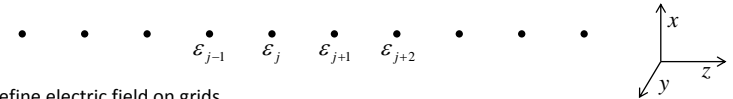
# Numerical solution

- Consider 1D Maxwell's equations

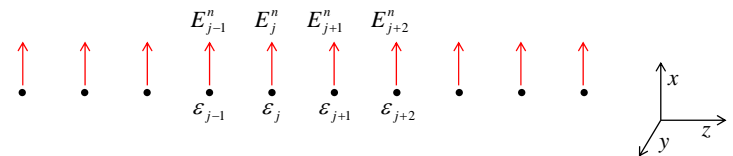
$$\frac{\partial E_x}{\partial z} = -\mu_0 \frac{\partial H_y}{\partial t}, \quad \frac{\partial H_y}{\partial z} = -\epsilon \frac{\partial E_x}{\partial t}$$

- Finite-difference time-domain method (FDTD)

- discretize space into grids ( $\Delta x$  distance between grids)



- Define electric field on grids



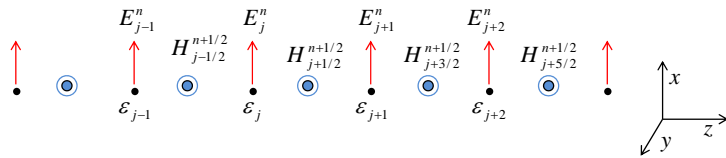
- we will also discretize time into time-steps ( $\Delta t$  time between time-steps)

$$E_j^{n-1} \quad E_j^n \quad E_j^{n+1} \quad E_j^{n+2}$$

- put magnetic field at half grids and half time-steps  $H_{j-1/2}^{n+1/2} \quad H_{j+1/2}^{n+1/2} \quad H_{j+3/2}^{n+1/2}$

## 1D FDTD

- Space-time grid



- Approximate Maxwell's equations by central differences

– Assume we know everything for  $E$  at time  $n-1$  and for  $H$  at time  $n-1/2$

$$\frac{\partial E_x}{\partial t} = -\frac{1}{\varepsilon} \frac{\partial H_y}{\partial z} \quad \rightarrow \quad \frac{E_j^n - E_j^{n-1}}{\Delta t} = -\frac{1}{\varepsilon_j} \frac{H_{j+1/2}^{n-1/2} - H_{j-1/2}^{n-1/2}}{\Delta z}$$

$$\frac{\partial H_y}{\partial t} = -\frac{1}{\mu_0} \frac{\partial E_x}{\partial z} \quad \rightarrow \quad \frac{H_{j+1/2}^{n+1/2} - H_{j+1/2}^{n-1/2}}{\Delta t} = -\frac{1}{\mu_0} \frac{E_{j+1}^n - E_{j-1}^n}{\Delta z}$$

- Field update algorithm

$$E_j^n = E_j^{n-1} - \frac{c\Delta t}{\Delta z} \frac{1}{\varepsilon_j} (H_{j+1/2}^{n-1/2} - H_{j-1/2}^{n-1/2})$$

$$H_{j+1/2}^{n+1/2} = H_{j+1/2}^{n-1/2} - \frac{c\Delta t}{\Delta z} (E_{j+1}^n - E_{j-1}^n)$$

the normalization was made

$$E \rightarrow \sqrt{\varepsilon_0 / \mu_0} E$$

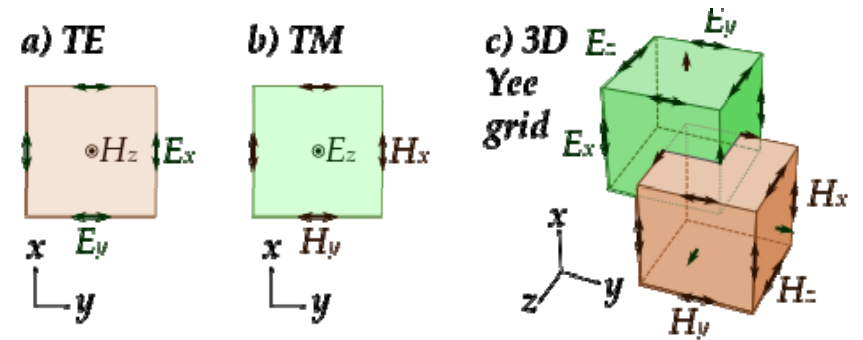
1D stability condition

$$c\Delta t < \Delta z$$

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## 2D and 3D FDTD

- Yee lattice: different sublattices for electric and magnetic fields
  - central differences for all field components



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